

Definite signature conformal holonomy: A complete classification

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Abstract

This paper aims to classify the holonomy of the conformal Tractor connection, and relate these holonomies to the geometry of the underlying manifold. The conformally Einstein case is dealt with through the construction of metric cones, whose Riemannian holonomy is the same as the Tractor holonomy of the underlying manifold. Direct calculations in the Ricci-flat case and an important decomposition theorem complete the classification for definitive signature.

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1. Introduction

Conformal geometry is perhaps the most natural extension of Riemannian geometry, and shares many of the same features with it. However, it was realised early on – as far back as Cartan [16] – that one of the most mathematically rewarding ways of dealing with conformal geometry was not by analogy with Riemannian geometries, but by analogy with the other parabolic geometries, using the general *Cartan connection* as a universal tool.

These parabolic geometries are a class of geometries that include, amongst others, projective, almost Grassmannian, almost quaternionic, and codimension one CR structures. The common point of these is that their ‘flat’ model space is the Lie group quotient G/P , where P is parabolic. Papers [1,2] by the same author deals with the projective case, while this paper treats the conformal one.

Many figures contributed to understanding parabolic geometries; Thomas [34,35] developed key ideas for Cartan connection calculus, and Shiego Sasaki investigated the conformal case in 1943 [32,33], followed by Tanaka [31] in 1979 and the major paper of Bailey, Eastwood and Gover in 1994 [4].

Since then, there have been a series of papers by Čap and Gover [14,13,20,15], developing a lot of the techniques that will be used in the present paper. Previous papers had focused on seeing the Cartan connection for conformal geometry as a property of a principal bundle \mathcal{P} . More recently, the principal bundle is replaced by an associated vector bundle, the *Tractor bundle* \mathcal{T} , and the Cartan connection by a connection form for \mathcal{T} , the *Tractor connection* $\overline{\nabla}$. With these tools, calculations are considerably simplified.

The purpose of this paper is to analyse one of the invariants of the Tractor connection, the holonomy group. There is an invariant metric of signature $(n + 1, 1)$ on \mathcal{T} , so this holonomy group must be a subgroup of $G = SO(n + 1, 1)$.

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It is a well known fact that a parallel section of the Tractor bundle corresponds to the local existence of an Einstein metric in the conformal class of a manifold. Beyond this, little was known about reductions of holonomy.

In this paper, we shall classify all the possible local holonomy groups of $\bar{\nabla}$ acting *reducibly* on \mathcal{T} . In doing so, they must conserve a Lorentzian metric of signature $(n + 1, 1)$. Then a paper by Di Scala and Olmos [17] shows that we have the complete list: there exist no connected proper subgroups of $SO(n + 1, 1)$ acting irreducibly on $\mathbb{R}^{n+1,1}$.

Proposition 1.1. *There are no local holonomy algebras acting irreducibly on \mathcal{T} apart from the full $\mathfrak{so}(n + 1, 1)$ algebra.*

A very recent paper by Felipe Leitner, [28], proves the same results as in this paper; but his methods, involving normal Killing forms, are different from those described here.

The classification comes in two main pieces; if a bundle of rank other than 1 or n is preserved, the manifold decomposes analogously to the De Rham decomposition:

Theorem 1.2. *Let $(M, [g])$ be a conformal, n -dimensional manifold, such that \mathcal{T}_M has a holonomy preserved subbundle of rank k , $2 \leq k \leq n$. Then locally there exists a metric $g \in [g]$ such that (M, g) splits locally into the direct product of two Einstein manifolds N_1, N_2 of dimensions $l = k - 1$ and $n - l$. The Einstein constants a and b of N_1, N_2 are related by $(n - l - 1)a = (1 - l)b$. Furthermore, there are canonical inclusions of the Tractor bundles of N_1 and N_2 into \mathcal{T}_M and the Tractor holonomy group of M is locally the direct product of those of N_1 and N_2 .*

That last statement requires a bit of explaining, since the Tractor bundles of N_1 and N_2 are of rank $l + 2$ and $n - l + 2$ respectively. However, since these are both Einstein manifolds, the effective rank of their Tractor bundles are $l + 1$ and $n - l + 1$, allowing the decomposition.

The converse to [Theorem 1.2](#) is also true. This decomposition is a local result, and may become degenerate along some embedded submanifolds.

The second step is to list all the possible Tractor holonomies for a conformally Einstein manifold. Using a metric cone construction, related to the ambient metric of [18,15,21], the following list is established:

Theorem 1.3 (Einstein Classification). *The Tractor holonomy of $(M^n, [g])$, for M^n conformal to an Einstein space of non-zero scalar curvature, is one of the following groups:*

- $SO(n, 1)$, $n \geq 4$,
- $SO(n + 1)$, $n \geq 4$,
- $SU(m)$ for $2m = n + 1$, $n \geq 4$,
- $Sp(m)$ for $4m = n + 1$,
- G_2 for $n = 6$,
- $Spin(7)$ for $n = 7$.

Moreover, all these actually occur as holonomy groups.

The Ricci-flat case must be treated differently; in fact, if (M^n, g) is Ricci-flat and conformally indecomposable, and G is the metric holonomy group of ∇^g , then $(M^n, [g])$ has Tractor holonomy $G \times \mathbb{R}^n$. Thus:

Theorem 1.4. *The possible indecomposable Tractor holonomy groups for the conformal manifold $(M^n, [g])$, conformally Ricci-flat, are:*

- $SO(n) \times \mathbb{R}^n$, $n \geq 4$,
- $SU(m) \times \mathbb{R}^{2m}$, $m \geq 2$,
- $Sp(m) \times \mathbb{R}^{4m}$, $m \geq 1$,
- $G_2 \times \mathbb{R}^7$,
- $Spin(7) \times \mathbb{R}^8$,

and all of these groups do occur.

This paper begins with defining and laying out the groundwork for the conformal Tractor bundle and connection. Furthermore, it will prove the equivalence of this (second-order) point of view with the standard view of the conformal structure as an equivalence class of metric structures. Some standard results will then be presented, showing how an Einstein structure in the conformal class is equivalent to a parallel section of the Tractor bundle.

Section 4, the heart of the paper, introduces umbilicity, the conformal equivalent of totally geodicity, and proves the decomposition theorem previously mentioned.

Section 5 then establishes the list for the Einstein spaces via the metric cone construction, with Section 6 complementing it using different methods to list the possible holonomies for conformally Ricci-flat manifolds.

A brief note on symmetric spaces follows, to illustrate the use of these methods; the paper ends with considerations of the differences that arise with indefinite signature.

This paper formed the beginning of the author's thesis [3] and was inspired and supervised by Dr. Nigel Hitchin.

Remark. In all the holonomy groups listed in this paper, the holonomy reduction corresponds to the existence of a particular metric in the conformal class. Hence we always have a canonical representative in the conformal class, whenever the holonomy reduces.

2. Cartan connection: Theory

2.1. The Cartan connection

With homogeneous geometries, since Klein, one deals with homogeneous spaces $M = G/P$, for G a Lie group acting transitively and effectively on X and P a subgroup.

This makes G into a principal P -bundle over M :

$$0 \rightarrow P \rightarrow G \rightarrow M \rightarrow 0.$$

The left-invariant vector fields on G give rise to an isomorphism $TG_u \rightarrow TG_e \cong \mathfrak{g}$ at every element $u \in G$. This isomorphism generates a 1-form ω on G , with values in the Lie algebra \mathfrak{g} .

The Cartan connection is a way of generalising these structures to a non-homogeneous M . We will follow the exposition used in [14]. Reading that article would explain most of the background and the reasons for the constructions detailed here.

In all of the following, we assume that M is an n -dimensional manifold, with \mathfrak{g} a semisimple Lie algebra and a subalgebra $\mathfrak{p} \subset \mathfrak{g}$ with \mathfrak{p} of codimension n in \mathfrak{g} . There are corresponding groups $P \subset G$; different choices of such groups may change the global properties of Cartan connections, but not the local ones.

Definition 2.1 (*Cartan Connection*). On M , given a principal P -bundle $\mathcal{P} \rightarrow M$, a Cartan connection ω is a section of $T_{\mathcal{P}}^* \otimes \mathfrak{g}$, with the following properties:

1. ω is invariant under the P -action (P acting by Ad on \mathfrak{g}),
2. $\omega(\sigma_A) = A$, where σ_A is the fundamental vector field of $A \in \mathfrak{p}$,
3. $\omega_u : T\mathcal{P}_u \rightarrow \mathfrak{g}$ is a linear isomorphism for all $u \in \mathcal{P}$.

If \mathfrak{p} is moreover a parabolic subalgebra (see paper [12]), we may make the further requirement that the connection be *normal*; this is a uniqueness condition for the Cartan connection of a particular geometry, similar to the torsion-free condition for a Levi-Civita connection. In general, this is the condition that the ‘curvature’ of the connection is closed under a certain Lie algebra homology, see [14]. There is, however, a simpler characterisation of the normality condition in the conformal case, see the proof of Lemma 2.14. See [14] for a proof of the existence of a normal Cartan connection in all parabolic geometries.

Paper [12] defines a parabolic subalgebra in an elegant and invariant way; for our purposes, however, it suffices to require that there exists a graded splitting of \mathfrak{g}

$$\mathfrak{g} = \mathfrak{g}_l \oplus \mathfrak{g}_{l-1} \oplus \cdots \oplus \mathfrak{g}_1 \oplus \mathfrak{g}_0 \oplus \mathfrak{g}_{-1} \oplus \cdots \oplus \mathfrak{g}_{-l},$$

such that $[\mathfrak{g}_j, \mathfrak{g}_k] \subseteq \mathfrak{g}_{j+k}$ and

$$\mathfrak{p} = \mathfrak{g}_0 \oplus \mathfrak{g}_{-1} \oplus \cdots \oplus \mathfrak{g}_{-l}.$$

The algebra is then called $|l|$ -graded; the conformal algebra will be seen to be $|1|$ -graded.

The bundle \mathcal{P} and the form ω together define the geometry. The first two conditions on ω are analogous to those of a standard principal bundle connection. The third condition is very different, however, giving a pointwise isomorphism $T\mathcal{P}_u \rightarrow \mathfrak{g}$ rather than a map with kernel.

However the Cartan connection does give rise to a connection in the usual sense, the so-called Tractor connection:

2.1.1. The tractor connection

The inclusion $P \hookrightarrow G$ generates a principal bundle inclusion $i : \mathcal{P} \hookrightarrow \mathcal{G}$, with \mathcal{G} a G -bundle, and generates a standard principal bundle connection form:

Proposition 2.2. *There is a unique $\omega' \in \Omega^1(\mathcal{G}, \mathfrak{g})$ such that ω' is a principal bundle connection form on \mathcal{G} and $i^*\omega' = \omega$.*

Proof. At any point of $\mathcal{P} \hookrightarrow \mathcal{G}$, define $\omega'(X) = \omega(X)$ for $X \in \Gamma(T\mathcal{P})$, and $\omega'(\sigma_A) = A$ for σ_A the fundamental vector field of $A \in \mathfrak{g}$. These two formulas correspond whenever they are both defined (Property 2 from Definition 2.1), and completely define ω' on \mathcal{P} . For any given $u \in \mathcal{G}$, let g be an element of G such that $g(u)$ is an element of \mathcal{P} . Then we may define ω' at u via the equality:

$$Ad_g \cdot \omega'_u = g^*(\omega'_{g(u)}),$$

Ad_g acting on \mathfrak{g} . Property 1 for ω ensures this is well defined.

To see that ω' is indeed a connection, notice that for $v \in \mathcal{P}$, $\omega' : T\mathcal{G}_v \rightarrow \mathfrak{g}$ has maximal rank, since $\omega = \omega'|_{T\mathcal{P}} : T\mathcal{P}_v \rightarrow \mathfrak{g}$ is surjective. G -invariance of ω' generalises this property to all of \mathcal{G} . \square

This ω' is the Tractor connection; when we see it as a connection on an associated vector bundle, we shall designate it by $\vec{\nabla}$. The Tractor connection obviously generates a Cartan connection by pull-back to $T\mathcal{P}$. From now on, we shall use Cartan and Tractor connections interchangeably.

Remark. It is not the case that any G connection η will correspond to a Cartan connection via pull-back to \mathcal{P} , as the isomorphism condition $T\mathcal{P}_v \rightarrow \mathfrak{g}$ could be violated. In fact, η must have a maximal second fundamental form on the canonical subbundles in the splitting of the Tractor bundle. This form is sometimes known as the soldering form [10]. If so, then η comes from a Cartan connection.

2.2. Conformal geometry

There are thus three standard ways of envisaging conformal geometry on a manifold M :

- via a class of conformal metrics $[g]$ related by multiplication by a never-zero function (a zero-order structure),
- via a class of torsion-free conformal connections ∇ (a first-order structure) or
- via a Cartan/Tractor connection $\omega/\vec{\nabla}$ (a second-order structure),

We will give more details of these three structures, and show their equivalence. The equivalence is easy to see between the first two structures – $[g]$ defines a conformal frame bundle which is the principal bundle for the connections ∇ – but is non-trivial with the third structure.

Let $\mathfrak{co}(n) = \mathfrak{so}(n) \oplus \mathbb{R}$ be the conformal algebra, with conformal group $CO(n)$. Then let \mathcal{G}_0 be the principal $CO(n)$ -frame bundle defined by $[g]$.

This allows us to define the bundles $\mathcal{E}[w]$, the weighted line bundles coming from the centre of $CO(n)$, i.e.

$$\mathcal{E}[w] = \mathcal{G}_0 \times_{\rho} \mathbb{R}, \quad \rho(c)(z) = -w \frac{\det(c)}{2n} z.$$

It is easy to see that $\mathcal{E}[-n] = \wedge^n T^*$ and $\mathcal{E}[w] = \mathcal{E}[-n]^{-\frac{w}{n}}$, so any connection on the tangent bundle extends to a connection on these weighted line bundles.

From now on, we will use the notation $B[w]$ for $B \otimes \mathcal{E}[w]$. Then there is a map from any $g \in [g]$ to a section of $(\odot^2 T)[-2]$,

$$g \rightarrow \mathbf{g} = \frac{n}{\det g} g.$$

The metric g determines a metric on $\otimes T^*$ by extending g via the relation $g(a \otimes b, a \otimes b) = g(a, a)g(b, b)$. Then $\det g$ is the section of $\mathcal{E}[-n] = \wedge^n T^* \subset \otimes T^*$ whose norm is one or minus one at every point of the manifold. Picking a g -orthonormal basis for T demonstrates that $\det g$ exists whenever g is non-degenerate.

Rescaling g demonstrates that \mathbf{g} does not depend on the choice of g . Conversely, given a non-vanishing section ξ of $\mathcal{E}[1]$ – a *conformal scale* – there is a corresponding metric in the conformal class

$$g^\xi = \xi^{-2}\mathbf{g},$$

with a corresponding Levi-Civita connection ∇ . Thus the class $[g]$ and the *conformal metric* \mathbf{g} are equivalent, and we will use them interchangeably.

The second way of defining the conformal structure is to use the class of preferred connections:

Definition 2.3. Given a conformal manifold $\mathcal{G}_0 \rightarrow M$, a *preferred connection* ∇ is a torsion-free \mathcal{G}_0 connection.

Proposition 2.4. Given a conformal structure, a preferred connection is equivalent with a connection on $\mathcal{E}[1]$ (or on any $\mathcal{E}[a] = (\mathcal{E}[1])^a$, for $a \neq 0$).

Proof. Similarly with the Levi-Civita connection expression, given \mathbf{g} and any connection ∂ on $\mathcal{E}[2]$, define ∇ for sections X, Y and Z of T as:

$$2\mathbf{g}(\nabla_X, Y) = \partial_X \mathbf{g}(Y, Z) + \partial_Y \mathbf{g}(Z, X) - \partial_Z \mathbf{g}(X, Y) + \mathbf{g}([X, Y], Z) - \mathbf{g}([X, Z], Y) - \mathbf{g}([Y, Z], X).$$

This ∇ is torsion-free. Conversely, each ∇ defines a connection on $\mathcal{E}[-n]$ and hence on any $\mathcal{E}[a]$, $a \neq 0$. This means that for a given conformal structure, ∂ and ∇ are in one-to-one correspondence, and we shall designate them both by ∇ . \square

In this view, those preferred connections that preserve a metric are exclusively those that preserve a conformal scale — and hence have trivial curvature on $\mathcal{E}[w]$.

2.2.1. The Cartan connection

In the classical, flat, case, conformal geometry is modelled on the sphere S^n . Taking the sphere as the collection of null-lines in $\mathbb{R}^{n+1,1}$, the group G of conformal transformations is $SO(n + 1, 1)$ in for n odd, and is double-covered by $SO(n + 1, 1)$ for even n . Then its Lie algebra has a 1-grading,

$$\mathfrak{g} = \mathbb{R}^n \oplus \mathfrak{co}(n) \oplus \mathbb{R}^{n*},$$

where the conformal group $\mathfrak{co}(n)$ decomposes into the semisimple part $\mathfrak{so}(n)$ and the centre \mathbb{R} , which is responsible for the conformal weight in representations of $\mathfrak{co}(n)$.

Thus the data are $\mathfrak{g} = \mathfrak{so}(n + 1, 1)$, $\mathfrak{g}_0 = \mathfrak{co}(n)$ and $\mathfrak{p} = \mathfrak{co}(n) \times \mathbb{R}^{n*}$, on an n -dimensional space. Note that we have a natural action of \mathfrak{g}_0 on \mathfrak{g} and hence an associated bundle to the \mathcal{G}_0 structure bundle:

$$\mathcal{G}_0 \times_{G_0} \mathfrak{g}.$$

Moreover, the action of G_0 splits \mathfrak{g} , giving a corresponding splitting:

$$\mathcal{G}_0 \times_{G_0} \mathfrak{g} = T \oplus \mathfrak{co}(T) \oplus T^*.$$

This decomposition will be used extensively.

It is important to explicate the Lie bracket of this algebra [14]. In fact, $[T, T] = [T^*, T^*] = 0$, the Lie bracket on $\mathfrak{co}(T)$ is the natural commutator of endomorphisms, and $[c, t^*] = -c(t^*)$, $[c, t] = c(t)$, for t, c, t^* sections of $T, \mathfrak{co}(T)$ and T^* respectively. The bracket between T and T^* is more complicated, and in fact

$$[t, s^*] = t \otimes s^* - (t \otimes s^*)^\tau + s^*(t)\delta,$$

with τ the transpose operator, and δ the identity element in (the centre of) $\mathfrak{co}(T)$.

In their papers [14,13], the authors demonstrate that the Cartan connection is equivalent to the standard conformal structure on a manifold $(M, [g])$. This is an alternative treatment.

Theorem 2.5. *Let s be a section of any bundle associated to \mathcal{G}_0 , and let X be any vector field. Then if ∇ and $\widehat{\nabla}$ are two preferred connections, there exists a 1-form \mathcal{Y} such that*

$$\widehat{\nabla}_X s = \nabla_X s + [\mathcal{Y}, X].s, \tag{1}$$

where $[\cdot, \cdot]$ is the Lie bracket for \mathfrak{g} previously described.

Proof. For $X \in \Gamma(T)$ and $\mathcal{Y} \in \Omega^1(M)$, $[\mathcal{Y}, X]$ is a section of $\mathfrak{co}(T)$, so this identity makes sense.

We know that $\widehat{\nabla}_X = \nabla_X + q(X)$, where q is a 1-form with values in $\mathfrak{co}(T)$.

However preferred connections are torsion-free, so $q(X).Y$ must be symmetric in X and Y , implying that q lies in the bundle $Q = (\odot^2 T^* \otimes T) \cap (T^* \otimes \mathfrak{co}(T))$, the symmetrisation of $T^* \otimes \mathfrak{co}(T)$ around the first two elements.

The fact that Q is of rank n and spanned by elements of the form $[\mathcal{Y}, -]$ can be seen by fairly simple Lie algebra manipulations (for more details, see [3]; the important idea is that $Q \cap T^* \otimes \mathfrak{so}(T) = 0$ by the uniqueness of the Levi-Civita connection, ensuring that the rank of Q is $\leq n$). \square

Note that if ∇ were a metric connection, then $\widehat{\nabla}$ would be metric if and only if \mathcal{Y} were a closed form. In fact:

Proposition 2.6. *Let ∇ and $\widehat{\nabla} = \nabla + \mathcal{Y}$ be two metric, preferred connections, with $\xi, \widehat{\xi}$ the corresponding conformal scales. Defining the function f as $f = \xi \otimes \widehat{\xi}^{-1} \in \Gamma(\mathcal{E}[0]) = C^\infty(M)$, we have*

$$\mathcal{Y} = f^{-1}df = d(\log f).$$

Proof. By direct calculation, using the fact that ∇ annihilates g^ξ while $\widehat{\nabla}$ annihilates $g^{\widehat{\xi}} = f^2(g^\xi)$. \square

A variety of tensors connected with these preferred connections will be needed in subsequent chapters. To define them, we will use Penrose’s abstract index notation, where Q^i is understood as a section of the tangent bundle, Q_i a section of the cotangent bundle, and symmetric and anti-symmetric parts over i and j to be denoted by (ij) and $[ij]$ respectively. This notation will be used intermittently throughout the paper.

Then if R_{ijkl} is the curvature tensor of ∇ , recall [20]:

$$R_{ijkl} = W_{ijkl} + 2\mathbf{g}_{k[i}P_{j]l} - 2\mathbf{g}_{l[i}P_{j]k} - 2P_{[ij]}\mathbf{g}_{kl} \tag{2}$$

with W_{ijkl} the conformally invariant Weyl tensor, and the rho tensor P :

$$P_{ij} = -\frac{1}{n-2} \left(\frac{1}{n} \text{Ric}_{ij} + \frac{n-1}{n} \text{Ric}_{ji} - \frac{1}{2n-2} R\mathbf{g}_{ij} \right) \tag{3}$$

a particularly important tensor for the rest of the paper. Here, Ric_{ij} is the Ricci curvature, and R the scalar curvature $\text{Ric}_{ij}\mathbf{g}^{ij}$ — a section of $\mathcal{E}[2]$.

This is in the general case for a conformal connection; in the metric case, the picture is the same, except that P follows the simpler symmetric formula

$$P_{ij} = -\frac{1}{n-2} \left(\text{Ric}_{ij} - \frac{1}{2n-2} R\mathbf{g}_{ij} \right).$$

The last relevant tensor for ∇ is the Cotton–York tensor:

$$CY_{ijk} = 2\nabla_{[i}P_{j]k}. \tag{4}$$

It will be important to understand how the tensor P varies under a change of conformal structure, as this formula is the key to defining the Tractor bundle. Letting P be the rho tensor for ∇ and \widehat{P} be that of $\widehat{\nabla} = \nabla + \mathcal{Y}$,

$$\widehat{P}(\xi) = P(\xi) - \nabla_\xi \mathcal{Y} + \frac{1}{2}[\mathcal{Y}, [\mathcal{Y}, \xi]], \tag{5}$$

for ξ any vector field.

2.2.2. Equivalences

Here we will demonstrate the equivalence of the Cartan connection with the conventional conformal structure. Though we will draw heavily on [13] for this exposition, we will use a slightly unconventional approach, which has the advantage of constructing the vital ‘Tractor bundle’ directly.

Remark. For a variety of reasons to do mainly with conventional notation and ease of calculations, we will be working with the Tractor bundle \mathcal{T} in the rest of the paper. However, to get a better understanding of what this bundle actually is, we need to start by defining the dual bundle \mathcal{T}^* .

Consider the two-jet prolongation of $J^2(\mathcal{E}[1])$ of the weighted bundle $\mathcal{E}[1]$. By definition, we have the short exact sequences

$$\begin{aligned} 0 &\longrightarrow \odot^2 T^*[1] \longrightarrow J^2(\mathcal{E}[1]) \longrightarrow J^1(\mathcal{E}[1]) \longrightarrow 0, \\ 0 &\longrightarrow T^*[1] \longrightarrow J^1(\mathcal{E}[1]) \longrightarrow \mathcal{E}[1] \longrightarrow 0. \end{aligned}$$

The conformal structure \mathbf{g} contracts $\odot^2 T^*$ to $\mathcal{E}[-2]$. Hence $\odot^2 T^*$ splits as $(\odot^2 T^*)_0 \oplus \mathcal{E}[-2]$, where the first space is the kernel of the contraction. Then we define the dual Tractor bundle \mathcal{T}^* as the quotient:

$$0 \longrightarrow (\odot^2 T^*)_0[1] \longrightarrow J^2(\mathcal{E}[1]) \longrightarrow \mathcal{T}^* \longrightarrow 0.$$

It is actually possible to realise \mathcal{T}^* as a subbundle of $J^2(\mathcal{E}[1])$ rather than a quotient bundle; we shall not be needing this result, though. Let D be the second-order operator $\Gamma(\mathcal{E}[1]) \rightarrow \Gamma(\mathcal{T}^*)$ given by composing the projection $J^2(\mathcal{E}[1]) \rightarrow \mathcal{T}^*$ with the two-jet operator j^2 .

Proposition 2.7. *Given a preferred connection ∇ , s any section of $\mathcal{E}[1]$ and b any point on the manifold, the map*

$$Ds(b) \rightarrow \left(s(b), \nabla_i s(b), \frac{1}{n} \mathbf{g}^{ij} (-\nabla_i \nabla_j s(b) + \mathbf{P}_{ij} s(b)) \right)$$

generates an isomorphism $\mathcal{T}^ \rightarrow \mathcal{E}[1] \oplus T^*[1] \oplus \mathcal{E}[-1]$.*

Proof. This formula clearly generates a bundle map $J^2(\mathcal{E}[1]) \rightarrow \mathcal{E}[1] \oplus T^*[1] \oplus \mathcal{E}[-1]$. All that remains is to prove that $(\odot^2 T^*)_0[1]$ is the kernel of this map. Assume $Ds(b) = 0$.

Then obviously $j^1(s) = 0$ at b , implying that $\nabla_i \nabla_j s(b)$ is the (well-defined) section of $\odot^2 T^*$ that corresponds to the second derivative of s at b . Thus $-\mathbf{g}^{ij} \nabla_i \nabla_j s(b) = 0$, or equivalently $j^2(s)(b) = \nabla_i \nabla_j s(b) \in (\odot^2 T^*)_0$. \square

Paper [13] directly demonstrates the following proposition:

Proposition 2.8. *Under a change of preferred connection $\nabla \rightarrow \widehat{\nabla} = \nabla + \Upsilon$, the isomorphism of Proposition 2.7 changes as*

$$\begin{pmatrix} x \\ \omega_i \\ z \end{pmatrix} \rightarrow \begin{pmatrix} x \\ \omega_i - \Upsilon_i x \\ z + \mathbf{g}^{ij} \omega_j \Upsilon_i - \frac{1}{2} \mathbf{g}^{ij} \Upsilon_i \Upsilon_j x \end{pmatrix}.$$

However, for reasons of convenience and notation, we will be working not with the bundle \mathcal{T}^* but with its dual. Define the *Tractor bundle* as $\mathcal{T} = (\mathcal{T}^*)^*$.

The previous results carry through to the dual of \mathcal{T}^* ; any preferred connection ∇ defines a splitting $\mathcal{T} = \mathcal{E}[1] \oplus T[-1] \oplus \mathcal{E}[-1]$, and under a change of connection, this splitting changes via

$$\begin{pmatrix} x \\ Y \\ z \end{pmatrix} \rightarrow \begin{pmatrix} x \\ Y + \Upsilon^* x \\ z - \Upsilon(Y) - \frac{1}{2} \mathbf{g}(\Upsilon, \Upsilon)x \end{pmatrix},$$

where $\Upsilon^* \in T[-2]$ is the dual to Υ using the conformal metric \mathbf{g} .

This particularly nice change of splitting formula implies the next vital lemma:

Lemma 2.9. *There is a natural metric $\langle \cdot, \cdot \rangle$, of type $(n + 1, 1)$, on \mathcal{T} .*

Proof. Given a preferred connection ∇ and two sections of \mathcal{T} , (x, Y, z) and (x', Y', z') , we define the metric by

$$\left\langle \begin{pmatrix} x \\ Y \\ z \end{pmatrix}, \begin{pmatrix} x' \\ Y' \\ z' \end{pmatrix} \right\rangle = xz' + x'z + \mathbf{g}(Y, Y').$$

Direct calculation then shows this formula is invariant under a change of splitting. \square

Since \mathcal{T}^* came about as a quotient bundle of a jet bundle, we have invariant subspaces of \mathcal{T} :

$$\mathcal{E}[-1] \hookrightarrow \mathcal{T} \times \mathcal{E}[-1] \hookrightarrow \mathcal{T},$$

and invariant projections:

$$\pi^1 : \mathcal{T} \rightarrow \mathcal{E}[1] \times \mathcal{T},$$

$$\pi^2 : \mathcal{T} \rightarrow \mathcal{E}[1].$$

Call E the subbundle of \mathcal{T} that is the inclusion of $\mathcal{E}[-1]$. Note that E is null under $\langle \cdot, \cdot \rangle$.

Lemma 2.10. *There is a P -bundle \mathcal{P} which is a principal bundle for \mathcal{T} , where P is the subgroup of G with Lie algebra \mathfrak{p} .*

Proof. The metric $\langle \cdot, \cdot \rangle$ shows that the structure algebra of \mathcal{T} reduces to $\mathfrak{g} = \mathfrak{so}(n + 1, 1)$. The invariant null subbundle E further reduces the structure algebra to $\{z \in \mathfrak{g} | z(E) = 0\}$, i.e. to \mathfrak{p} .

Then we define \mathcal{P} to be the bundle of orthonormal frames of \mathcal{T} preserving E . \square

Let us review what we have so far. Starting from the conformal metric \mathbf{g} and the class of preferred connections, we have constructed, via a tensor \mathbf{P} dependent on the connections, a bundle \mathcal{T} . And this bundle generates a principal bundle \mathcal{P} , where it is natural to suppose the Cartan connection living. We now need to build this Cartan connection.

Define the Lie algebra bundle $\mathcal{A} = \mathcal{P} \times_P \mathfrak{g}$. Then given a preferred connection ∇ we have a splitting of $\mathcal{T} = \mathcal{E}[1] \oplus T[-1] \oplus \mathcal{E}[-1]$, and hence a corresponding splitting:

$$\mathcal{A} = T \oplus \mathfrak{co}(T) \oplus T^*.$$

In order to finish the construction of this Cartan connection, we will start by building a $G = SO(n + 1, 1)$ connection on \mathcal{A} and then prove that it is a Tractor connection.

Definition 2.11. Given a preferred connection ∇ , we have a splitting of $\mathcal{A} = T \oplus \mathfrak{co}(T) \oplus T^*$. Each of these bundles is a G_0 -bundle, so ∇ ascends to a connection on \mathcal{A} . Then we define the Tractor connection $\overrightarrow{\nabla}$ as

$$\overrightarrow{\nabla}_X = \nabla_X + \text{ad}(X) + \text{adP}(X),$$

with the vector X and the 1-form $\mathbf{P}(X)$ seen as sections of the Lie algebra bundle \mathcal{A} .

Then since ∇ comes from a connection on a principal bundle with structure group G_0 , $\overrightarrow{\nabla}$ comes from a principal bundle connection with structure group G . Of course, this definition makes no sense without:

Proposition 2.12. *This definition is independent of the choice of ∇ .*

Proof. The formula for the change of splitting of \mathcal{A} (deduced directly from that of \mathcal{T}) is:

$$\begin{pmatrix} X \\ \Psi \\ \omega \end{pmatrix} \rightarrow \begin{pmatrix} X \\ \Psi + [\mathcal{T}, X] \\ \omega + [\mathcal{T}, \Psi] + \frac{1}{2}[\mathcal{T}, [\mathcal{T}, X]] \end{pmatrix}.$$

Then a direct calculation proves the result. \square

Thus for any bundle B associated to \mathcal{G} , we have an invariant connection form:

$$\vec{\nabla} = \nabla + \rho(X) + \rho\mathbf{P}(X).$$

In the case of the Tractor bundle \mathcal{T} , the detailed expression is:

$$\vec{\nabla}_X \begin{pmatrix} x \\ Y \\ z \end{pmatrix} = \begin{pmatrix} \nabla_X x - \mathbf{g}(X, Y) \\ \nabla_X Y + zX - x\mathbf{P}(X) \\ \nabla_X z + \mathbf{P}(X, Y) \end{pmatrix}. \tag{6}$$

Now we get to the result that ties all the structures together:

Theorem 2.13. *The connection $\vec{\nabla}$ is a normal Tractor connection.*

Proof. Let i be the inclusion $i : \mathcal{P} \hookrightarrow \mathcal{G}$, π projection $\pi : \mathcal{P} \rightarrow M$, and let $\omega \in \Omega^1(\mathcal{G}, \mathfrak{g})$ be the 1-form associated with the connection $\vec{\nabla}$.

We need to prove that $\mu = i^*(\omega)$ is an isomorphism $T\mathcal{P}_u \rightarrow \mathfrak{g}$ for all points $u \in \mathcal{P}$; then μ will be the Cartan connection generating the Tractor connection ω .

So now assume that μ is *not* an isomorphism at some point u , so there exists a vector $\xi \in T\mathcal{P}_u$ such that $\mu(\xi) = 0$. As vertical vectors in \mathcal{P} are mapped isomorphically onto \mathfrak{p} , $X = \pi_*(\xi)$ is a non-zero vector in $T_{\pi(u)}$.

Then define a local section j of M in \mathcal{P} , such that $j_*X = \xi$. This also gives us a section $i \circ j$ of \mathcal{G} . Then in the local frame bundle determined by this section, the connection $\vec{\nabla}$ is of the form

$$d + j^*i^*\omega = d + j^*\mu.$$

This shows that the bundle \mathcal{P} is infinitesimally conserved at $\pi(u)$ in the X direction, or, switching to the associated bundle $E = \mathcal{P} \times_{\mathfrak{p}} e$, that there is a section s of E , non-zero at $\pi(u)$, such that $\vec{\nabla}_X s = 0$ at $\pi(u)$.

However, the connection on the Tractor bundle is given by Eq. (6):

$$\vec{\nabla}_X \begin{pmatrix} 0 \\ 0 \\ z \end{pmatrix} = \begin{pmatrix} 0 \\ zX \\ \nabla_X z \end{pmatrix},$$

which is a contradiction as s (hence z) is non-zero at $\pi(u)$. So $\vec{\nabla}$ is indeed a Tractor connection. \square

And finally, to complete the circle:

Lemma 2.14. *The Cartan connection generated by $\vec{\nabla}$ is normal.*

Proof. By [14,26], this result is equivalent with the curvature of $\vec{\nabla}$ lying in the Lie algebra bundle of $\mathcal{G} \times_G \mathfrak{p}$. Alternately, the curvature must preserve the canonical bundle E .

In abstract index notation, the expression for $\vec{\nabla}_i \vec{\nabla}_j$ is:

$$\begin{aligned} \vec{\nabla}_i \vec{\nabla}_j &= \nabla_i \nabla_j + (\delta_j^l \nabla_i + \delta_i^l \nabla_j) + (\mathbf{P}_{jk} \nabla_i + \mathbf{P}_{ik} \nabla_j) + \nabla_i (\mathbf{P}_{jk}) \\ &\quad + \delta_i^l \circ \delta_j^m + \delta_i^l \circ \mathbf{P}_{jn} + \mathbf{P}_{ik} \circ \delta_j^m + \mathbf{P}_{ik} \circ \mathbf{P}_{jn}. \end{aligned}$$

This whole expression is a section of $(\otimes^2 T^*) \otimes (\otimes^2 \mathcal{A})$. Recall that $\mathcal{A} = T \oplus \mathfrak{co}(T) \oplus T^*$, which explains why the various tensors here have differing rank. To get this expression we have used the connection ∇ on T to define the second covariant derivative; however, we could have used any other connection, as we are about to anti-symmetrise i and j . Upon doing this, the terms in brackets vanish. Moreover, $\rho(X) \circ \rho(Y) = \rho(Y) \circ \rho(X)$ and similarly for 1-forms, meaning that:

$$R_{ij}^{\vec{\nabla}} = R_{ij}^{\nabla} + 2(\nabla_{[i} (\mathbf{P}_{j]k}) + \delta_{[i}^l \circ \mathbf{P}_{j]n} + \mathbf{P}_{k[i} \circ \delta_{j]}^m)$$

(recall that $[ij]$ means taking the anti-symmetric part over the i and j components). Looking back at equations Eqs. (2) and (4), we see that this expression is the sum of the Weyl tensor and the Cotton–York tensor. Or, expressed in

more conventional notation, in the splitting of \mathcal{A} determined by ∇ :

$$R_{X,Y}^{\overrightarrow{\nabla}} = \begin{pmatrix} 0 \\ W(X, Y) \\ CY(X, Y) \end{pmatrix}.$$

Since only $T \subset \mathcal{A}$ has a non-trivial action on the canonical bundle $E \subset \mathcal{T}$, this curvature expression must preserve E . \square

To construct the conformal structure from the Cartan connection is much simpler; indeed, the metric \langle, \rangle descends to the conformal metric \mathbf{g} on $T = E^\perp/E$.

2.2.3. One and two dimensions

Though any 2-manifold is conformally flat, with an infinite dimensional local conformal transformation group, paper [11] and other unpublished papers by the same author extend the concept of conformal Cartan connections to one and two dimensions, by constructing Möbius structures. As in higher dimensions, a choice of Weyl structure determines a splitting of the associated Tractor bundle. There is an ambiguity, however, in the trace-free symmetric part of the \mathbf{P} tensor; this may be chosen freely.

In this paper, the only case where Möbius structures will be needed are during the decomposition results of Section 4. In that decomposition, we will have a specific metric connection ∇ to work with on our submanifolds. For the results to be consistent with higher dimensional results, we make the assumption that the Möbius structure is given by

$$\mathbf{P}_{hj} = -\frac{1}{2}\text{Ric}_{jh},$$

in two dimensions, Ric the Ricci tensor of ∇ , and $\mathbf{P}_{hj} = 0$ in one dimensions.

3. Conformally Einstein manifolds

3.1. Important note

In most of the proofs in the remainder of this paper, it will be assumed that for a certain holonomy preserved subbundle $U \subset \mathcal{T}$ used in the proof, one has $\pi^2(U) \neq 0$. This will not be the case everywhere, of course; however:

Proposition 3.1. *Let $U \subset \mathcal{T}$ be a preserved subbundle under $\overrightarrow{\nabla}$. Then $\pi^2(U) \neq 0$ on Σ , an open, dense subset of M .*

Proof. Σ is open because of the $\pi^2(U) \neq 0$ condition.

Let $b \in M \setminus \Sigma$, and $u(b) = (0, Y(b), z(b))$ be a non-zero element of U_b . Then extend $u(b)$ locally to a section $u = (x, Y, z)$ of U by parallel transport along ‘rays’ from b . This implies that $\overrightarrow{\nabla}u = 0$ at b . Then picking any nowhere-zero section τ of $\mathcal{E}[-1]$, we can define the function $f : M \rightarrow \mathbb{R}$ by

$$f(c) = \pi^2(u)\tau.$$

The derivative of f is $(\nabla x)\tau + x(\nabla\tau)$. At b , this is just $(\nabla x)\tau$, and, since $\overrightarrow{\nabla}u = 0$ at b :

$$X.f(b) = (\nabla_X x)\tau = \mathbf{g}(X, Y)\tau.$$

If $Y \neq 0$ at b , this is non-zero for some X , so f is non-zero arbitrarily close to b . If $Y = 0$, then the first derivative is zero, and the second derivative is thus:

$$\begin{aligned} Z.(X.f)(b) &= (\nabla_Z \nabla_X x)\tau \\ &= -\mathbf{g}(X, \nabla_Z Y)\tau \\ &= -\mathbf{g}(X, -zZ)\tau. \end{aligned}$$

with $z(b) \neq 0$ as $u(b) \neq 0$. Then the second derivative is non-zero for $X = Z \neq 0$, for instance, forcing f to be non-zero arbitrarily close to b .

This implies that $\pi^2(u) \neq 0$ arbitrarily close to b , proving the result. \square

In fact, if the first derivative vanishes, b must be an *isolated* point.

The classic examples of this are the various conformally Einstein metrics on the sphere S^n . The sphere is conformally flat, so there are many holonomy preserved sections of its Tractor bundle.

A preserved section u of negative norm corresponds to the Spherical metric $g = \pi^2(u)^{-2}\mathbf{g}$ on the whole space. In this case, $\pi^2(u)$ is never zero.

A preserved section u of zero norm corresponds to the Euclidean metric $g = \pi^2(u)^{-2}\mathbf{g}$ on $\mathbb{R}^n \cong S^n \setminus \{\infty\}$. In this case, $\pi^2(u)(b) \neq 0$ for $b \neq \infty$.

A preserved section u of positive norm corresponds to the hyperbolic metric $g = \pi^2(u)^{-2}\mathbf{g}$ on two half-spheres of S^n . In this case $\pi^2(u)$ is zero only on the boundary S^{n-1} cutting S^n into two.

3.2. Einstein spaces

Though it is well known in general that any conformally Einstein space corresponds to a parallel section of the Tractor bundle \mathcal{T} , what follows is a direct proof of this fact using the Tractor connection approach. Paper [22] has a similar approach.

Remark. This is a first instance of a holonomy reduction of $\vec{\nabla}$.

Theorem 3.2. For $n > 2$, if (M, \mathbf{g}) has an Einstein metric g in its conformal class then there exists a parallel section s of its Tractor bundle \mathcal{T} .

Proof. Let g be the Einstein metric, $\text{Ric}^g = \lambda g$. Then the P tensor is

$$P = -\frac{\lambda}{2n-2}\mathbf{g}.$$

We may use the section ξ of $\mathcal{E}[1]$ corresponding to g to set up an isomorphism from functions $f : M \rightarrow \mathbb{R}$ to sections of $\mathcal{E}[a]$, by mapping f to $f\xi^a$. In this set up:

$$\vec{\nabla} \begin{pmatrix} 1 \\ 0 \\ \lambda \\ -\frac{\lambda}{2n-2} \end{pmatrix} = \begin{pmatrix} 0 \\ \frac{\lambda}{2n-2}Id - \frac{\lambda}{2n-2}Id \\ 0 \end{pmatrix} = 0. \quad \square$$

To prove the converse of this theorem, we need the following lemma:

Lemma 3.3. If a conformal connection ∇ has a symmetric Ricci tensor, then ∇ is actually a metric connection.

Proof. Let R^i_{ijkl} be the curvature of ∇ . Then R^i_{ijkl} acts on the determinant bundle $\mathcal{E}[n]$ via its trace R^i_{ikl} . However, by the first Bianchi identity,

$$\begin{aligned} R^i_{iikl} &= -R^i_{ikli} - R^i_{ilik} \\ &= \text{Ric}_{kl} - \text{Ric}_{lk}, \end{aligned}$$

the anti-symmetric part of the Ricci tensor. So if ∇ has a symmetric Ricci tensor, its curvature must vanish on $\mathcal{E}[n]$, so locally ∇ must preserve a section η of the determinant bundle. Then $\eta^{1/n}$ is a preserved conformal scale and

$$g = \eta^{-2/n}\mathbf{g}$$

is a metric preserved by that ∇ . \square

Theorem 3.4. For $n > 2$, if a line bundle L of \mathcal{T} is holonomy preserved, then a section s of L is preserved, and $(M, [g])$ has an Einstein metric $g = (\pi^2(s))^{-2}\mathbf{g}$ in its conformal class, wherever $\pi^2(s) \neq 0$.

Proof. The line bundle L defines a connection on $\mathcal{E}[1]$, and hence a torsion-free connection on T , in the following way. Let e be any nowhere-vanishing section of $\mathcal{E}[1]$, and let l be the section of L such that $\pi^2(l) = e$. Then define $\nabla e = \pi^2(\vec{\nabla}l)$; it is easy to see that this is indeed a connection.

Using ∇ , we split $\mathcal{T} = \mathcal{E}[1] \oplus T[-1] \oplus \mathcal{E}[-1]$ in the usual way. Then Eq. (6) implies that $\pi^2(\vec{\nabla}(x, Y, z)) = \nabla x - \mathbf{g}(Y, -)$. Since by definition of $\pi^2(\vec{\nabla}l) = \nabla\pi^2(l)$, we must have $Y = 0$ for l (and hence for any section of L).

If L is not null, then a section $s = (x, 0, x\mu)$ of constant norm, is preserved. This generates a metric $g = (\pi^2(s))^{-2}g$. But $\tilde{\nabla}_X(x, 0, x\mu) = (\nabla_X x, x\mu(X) - xP(X), \mu\nabla_X x) = 0$. This implies that $P = \mu g$, so $\text{Ric}^g = \lambda g$ for $\lambda = (2 - 2n)\mu$.

On the other hand, if L is null, $z = 0$, and $\tilde{\nabla}_X(x, 0, 0) = (\nabla_X x, -xP(X), 0)$. Thus $P = 0$ and hence ∇ has a symmetric Ricci tensor, implying that it is actually a metric connection for some metric g — which moreover is Ricci-flat. Set $x \in \Gamma(\mathcal{E}[1])$ to be the conformal scale corresponding to g . Then the section $s = (x, 0, 0)$, is parallel. \square

Proposition 3.5. *In the two-dimensional case, one merely has the one-way implication that an Einstein metric of constant scalar curvature gives a preserved section of \mathcal{T} .*

Remark. Note that the sign of $\langle s, s \rangle$ is the opposite of the sign of the Einstein constant λ .

4. Decomposition theorem

This section presents the decomposition theorem for Tractor connections, similar to the De Rham decomposition for Riemannian connections.

Remark. Related terminology may be found in [25].

4.1. Preparatory results

Definition 4.1. Given a metric g on M with Levi-Civita connection ∇ , a subbundle $U \subset T$ is *umbilical* for the connection ∇ , if there exists a vector field H such that for X and Y sections of U ,

$$\nabla_X Y = \tilde{\nabla}_X Y + g(X, Y)H,$$

for $\tilde{\nabla}$ some connection on U , and H a vector field.

Remark. Note that an umbilical subbundle is automatically integrable, as

$$\begin{aligned} [X, Y] &= \nabla_X Y - \nabla_Y X = \tilde{\nabla}_X Y - \tilde{\nabla}_Y X + (g(X, Y) - g(Y, X))H \\ &= \tilde{\nabla}_X Y - \tilde{\nabla}_Y X, \end{aligned}$$

a section of U .

Lemma 4.2. *U being umbilical is equivalent to*

$$\nabla_X Y \in \Gamma(U), \tag{7}$$

whenever X and Y are orthogonal sections of U .

Proof. If U is umbilical, then Eq. (7) is true by definition

$$\begin{aligned} \nabla_X Y &= \tilde{\nabla}_X Y + g(X, Y)H \\ &= \tilde{\nabla}_X Y \in \Gamma(U). \end{aligned}$$

So we now assume Eq. (7) and aim to prove umbilicity. One may easily see, by choosing an orthogonal frame for U , that U must be integrable.

Define a connection $\tilde{\nabla}$ on U , by orthogonal projection. Then the map $\Phi = \nabla - \tilde{\nabla}$ is bilinear, $U \otimes U \rightarrow U^\perp$, and symmetric since $\nabla_X Y - \nabla_Y X = [X, Y]$ is a section of U . By assumption, $\Phi(X, Y) = 0$ whenever $g(X, Y) = 0$.

Now let (X_j) be a frame of U , chosen so that the $g(X_j, X_k)$ are nowhere zero (one can do this, for instance, by choosing a standard orthonormal frame (X_j) and mapping $X_j \rightarrow X_j + \frac{1}{2n} \sum_{l=1}^j X_l$). Pick H in U^\perp such that $\Phi(X_1, X_1) = g(X_1, X_1)H$. Then since X_1 is orthogonal to $\tau_{1,1,j} = g(X_1, X_1)X_j - g(X_j, X_1)X_1$, one has $\Phi(X_1, \tau) = 0$ and hence

$$\begin{aligned} \Phi(X_1, X_j) &= \frac{1}{g(X_1, X_1)} g(X_j, X_1) (g(X_1, X_1)H) \\ &= g(X_j, X_1)H. \end{aligned}$$

The same argument with the orthogonal sections $\tau_{j,1,k}$ and X_j demonstrates

$$\Phi(X_j, X_k) = g(X_j, X_k)H.$$

This extends trivially to the whole of U . Thus $\nabla_X Y = \tilde{\nabla}_X Y + \Phi(X, Y) = \tilde{\nabla}_X Y + g(X, Y)H$. \square

Note that being umbilical is a conformally invariant condition, as changing ∇ by \mathcal{Y} changes H to $H - \mathcal{Y}^*$, since all the other terms in $[\mathcal{Y}, X] \cdot Y$ are tangent to U (see Eq. (1)). Thus choosing $\mathcal{Y}^* = H$, we can make U into a totally geodesic foliation. In other words, there are preferred connections for which U is totally geodesic.

4.2. Preserved subbundles

Let K be a subbundle of \mathcal{T} of rank k , $2 \leq k \leq n$, preserved by $\overrightarrow{\nabla}$. Then K defines a subbundle U of T as follows. We assume, from Proposition 3.1, that K and K^\perp are locally transverse to E . Recall that $E \cong \mathcal{E}[-1] \subset \mathcal{T}$ is the canonical line bundle, and that $E^\perp \cong TM[-1] \oplus L^{-1}$ is of rank $n + 1$ in \mathcal{T} .

Hence, $K \cap E^\perp$ is a bundle of rank $k - 1$, and π^1 is injective on $(K \cap E^\perp)$ (since $K \cap E = 0$, so π^1 is injective on K). Moreover $\pi^1(E^\perp) = T[-1]$, so

$$U = \pi^1(K \cap E^\perp) \subset T[-1]$$

is a well-defined, rank $k - 1$ bundle. Use any conformal scale to get an isomorphism $T \cong T[-1]$. Since changing the section simply results in scaling any element of $T[-1]$, we may see U as a well-defined subbundle of T .

Proposition 4.3. *U is an integrable, umbilical foliation of T .*

Proof. Let X and Y be orthogonal sections of U . Fix any metric in the conformal class. Then

$$\begin{pmatrix} 0 \\ Y \\ z \end{pmatrix},$$

is a section of $K \cap E^\perp$, for some z . Consequently

$$\overrightarrow{\nabla}_X \begin{pmatrix} 0 \\ Y \\ z \end{pmatrix} = \begin{pmatrix} 0 \\ \nabla_X Y + zX \\ z' \end{pmatrix},$$

for some z' . Since K is preserved by $\overrightarrow{\nabla}$, this is a section of K ; it is clearly a section of E^\perp . As a consequence, we know that

$$\nabla_X Y + zX \in \Gamma(U).$$

Thus $\nabla_X Y$ is also a section of $\Gamma(U)$, making U umbilical, and hence integrable. \square

We shall see later that U is Einstein (i.e. all leaves N of U are conformally Einstein under the restricted conformal structure).

Proposition 4.4. *There is a Tractor bundle \mathcal{T}_U on the leaves N of the foliation defined by U , and a well-defined inclusion $\mathcal{T}_U \subset \mathcal{T}$.*

Proof. If ∇ is a U -preferred connection – one that makes U , and its foliation, totally geodesic – in the splitting of \mathcal{T} that it defines,

$$\mathcal{T} = L^1 \oplus T[-1] \oplus L^{-1}.$$

Define \mathcal{T}_U as the subbundle

$$\mathcal{T}_U = L^1 \oplus U[-1] \oplus L^{-1}.$$

To check that this is well defined, we change ∇ to ∇' , another U -preferred connection. This is equivalent to changing ∇ by an $\Upsilon \in \Gamma(g(U) \subset T^*)$ for any metric g in the conformal class. Then the splitting changes as:

$$\begin{pmatrix} x \\ Y \\ z \end{pmatrix} \rightarrow \begin{pmatrix} x \\ Y + \Upsilon^*x \\ z - \Upsilon(Y) - \frac{1}{2}\mathbf{g}(\Upsilon, \Upsilon)x \end{pmatrix},$$

which, since Υ^* is a section of U , does not change the definition of \mathcal{T}_U nor its inclusion into \mathcal{T} . \square

We are now ready to prove the main theorem.

Theorem 4.5. *Assume there is a bundle K of rank k preserved by $\overrightarrow{\nabla}$, and the foliation U that it generates splits T . Let $l = k - 1$ be the rank of U . Then there exists a metric g in the conformal class of M such that the manifold (M, g) splits locally as the direct product*

$$(M, g) = (N_1, h_1) \times (N_2, h_2)$$

where h_1 and h_2 are Einstein metrics with Einstein coefficients λ_1, λ_2 , possibly zero, related by

$$(n - l - 1)\lambda_1 = (1 - l)\lambda_2.$$

The converse is also true. And in this situation the holonomy $\overrightarrow{\text{hol}}$ of $\overrightarrow{\nabla}$ is the direct sum of Lie algebras

$$\overrightarrow{\text{hol}} = \overrightarrow{\text{hol}}_{N_1} \oplus \overrightarrow{\text{hol}}_{N_2}$$

where $\overrightarrow{\text{hol}}_{N_1}$ is the holonomy of $\overrightarrow{\nabla}_{N_1}$ and $\overrightarrow{\text{hol}}_{N_2}$ that of $\overrightarrow{\nabla}_{N_2}$.

Note that the subbundle of T generated by K^\perp is just U^\perp . There are really two situations here: the case when $K \cap K^\perp$ is of rank 1, and that where it is of rank 0.

4.2.1. K degenerate

If $K \cap K^\perp = \mathcal{L}$, a line bundle, necessarily null, then by [Theorem 3.4](#) there must be a preserved section v of \mathcal{L} and hence a Ricci-flat metric g on M , with Levi-Civita connection ∇ .

We have the bundles U and U^\perp as before, both integrable and umbilical. We will now show that g is locally a product metric of the leaves generated by U and U^\perp . First, we shall demonstrate that these leaves are totally geodesic under g .

Lemma 4.6. *Let X be a section of U . Then for any $A \in \Gamma(T)$, $\nabla_A X$ is a section of U .*

Proof. In the splitting defined by g , one section of K is the Einstein vector

$$v = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}.$$

Since v is also a section of K^\perp , K must lie in v^\perp . In other words, K is of the form

$$\begin{pmatrix} \mathbb{R} \\ U \\ 0 \end{pmatrix}.$$

Now consider

$$\overrightarrow{\nabla}_A \begin{pmatrix} 0 \\ X \\ 0 \end{pmatrix} = \begin{pmatrix} -\mathbf{g}(A, X) \\ \nabla_A X \\ 0 \end{pmatrix}.$$

Since $\overrightarrow{\nabla}$ preserves K , $\nabla_A X$ must be a section of U . \square

This shows that U (and U^\perp) are totally geodesic foliations. Moreover, they are preserved by ∇ in every direction.

Remark. As a consequence of that, if X and B are commuting sections of U and U^\perp respectively,

$$\nabla_X B = \nabla_B X = 0.$$

Let $h_1 = g|_U$, Y and X be sections of U , A any section of T . Then

$$\begin{aligned} (\nabla_A h_1)(X, Y) &= A.h_1(X, Y) - h_1(\nabla_A X, Y) - h_1(X, \nabla_A Y) \\ &= A.g(X, Y) - g(\nabla_A X, Y) - g(X, \nabla_A Y) \\ &= (\nabla_A g)(X, Y) \\ &= 0, \end{aligned}$$

as $\nabla_A X$ and $\nabla_A Y$ are sections of U , and $h_1 = g$ on sections of U . Consequently we have demonstrated, for h_1 and for $h_2 = g|_{U^\perp}$:

Lemma 4.7. ∇h_1 and ∇h_2 are both zero.

Now pick sections X and Y of U commuting with a section B of U^\perp . By the previous lemma

$$B.h_1(X, Y) = 0,$$

so the Lie derivative of h_1 in the direction of B is

$$(\mathcal{L}_B h_1)(X, Y) = B.h_1(X, Y) - h_1([B, X], Y) - h_1(X, [B, Y]) = 0.$$

We may choose local coordinates that respect the foliations U and U^\perp to get frames (X^j) of U and (B^k) of U^\perp , commuting with one another. Consequently, if N_1 is a leaf of U and N_2 a leaf of U^\perp , h_1 is preserved by translation along N_2 and vice versa. This demonstrates that

Proposition 4.8. *Locally, $(M, g) = (N_1, h_1) \times (N_2, h_2)$.*

This implies that $\nabla|_U$ is the Levi-Civita connection of h_1 , and $\nabla|_{U^\perp}$ that of h_2 . To finish this exploration, we require:

Lemma 4.9 (Restricted Ricci Curvature). *Given a foliation U preserved by ∇ , the Ricci tensor of $\nabla|_U$ is the Ricci tensor of ∇ , restricted to U .*

Proof. Notice that this condition makes U integrable and totally geodesic. Let $(X_j), (B_j)$ be a coordinate frame for T , with $X_j \in \Gamma(U)$ and the (B_j) complementary. Then

$$\text{Ric}(X_j, X_k) = \left(\sum_l X_l^* \lrcorner R_{X_l, X_j} X_k \right) + \left(\sum_l B_l^* \lrcorner R_{B_l, X_j} X_k \right).$$

But the second term on the right is zero, as $R_{-, -} X_j$ must be a section of U , and the first term is just the Ricci curvature of ∇_U . \square

Consequently, one can see that ∇ is Ricci-flat on U and on U^\perp (hence on N_1 and N_2).

In this case the relation

$$(n - l - 1)\lambda_1 = (1 - l)\lambda_2$$

is trivially satisfied, as both λ_j are zero.

The converse to this construction is trivial: a direct product of Ricci-flat spaces is Ricci-flat. Then K may be reconstructed as

$$K = \begin{pmatrix} \mathbb{R} \\ TN_1 \\ 0 \end{pmatrix}$$

in the global Ricci-flat metric’s splitting. Since TN_1 must be totally geodesic, $\overrightarrow{\nabla}$ preserves K and

$$K^\perp = \begin{pmatrix} \mathbb{R} \\ TN_2 \\ 0 \end{pmatrix}.$$

Now notice that since all P are zero, $\overrightarrow{\nabla}$ acts on \mathcal{T}_{N_1} along N_1 exactly as the Tractor connection $\overrightarrow{\nabla}_{N_1}$ does. Moreover, $\overrightarrow{\nabla}$ acts trivially on \mathcal{T}_{N_1} along N_2 . Since the opposite result holds for \mathcal{T}_{N_2} , and since these two Tractor bundles span all of \mathcal{T} , one has

$$\overrightarrow{h}\sigma\overrightarrow{l} = \overrightarrow{h}\sigma\overrightarrow{l}_{N_1} \oplus \overrightarrow{h}\sigma\overrightarrow{l}_{N_2}.$$

4.2.2. K non-degenerate

We seek to imitate the proofs of the previous section in the case where $K \cap K^\perp = 0$. First of all, we seek to find an imitation of the Ricci-flat metric g . We shall use a preferred connection rather than a metric — though it will turn out to be a metric connection in the end.

Starting off, pick ∇' such that U is totally geodesic. In the rest of these proofs, X and Y will be sections of U , B and C sections of U^\perp .

Since U^\perp is umbilical,

$$\nabla'_B C = \tilde{\nabla}'_B C + H\tilde{g}(B, C),$$

for some $H \in \Gamma(U)$ and any metric \tilde{g} in the conformal class. Then replace ∇' with ∇ , by adding the 1-form $\Upsilon = \tilde{g}(H)$. This connection makes U^\perp totally geodesic, but since

$$\nabla_X Y = \nabla'_X Y + \Upsilon(X)Y + \Upsilon(Y)X - H\tilde{g}(X, Y)$$

is a section of U , then the bundle U remains totally geodesic under ∇ . In fact ∇ is the sole preferred connection that makes U and U^\perp totally geodesic — as adding any $\Upsilon \neq 0$ would destroy this property on at least one of these bundles.

Now we try and calculate K and K^\perp in the splitting given by ∇ . We know that elements of $K \cap E^\perp$ are of the form

$$\begin{pmatrix} 0 \\ X \\ z \end{pmatrix},$$

for some $z \in \Gamma(L^{-1})$ depending on X , and elements of $K \cap E^\perp$ are of the form

$$\begin{pmatrix} 0 \\ B \\ z' \end{pmatrix}.$$

Hence, choosing Y such that Y and X are not orthogonal,

$$\overrightarrow{\nabla}_Y \begin{pmatrix} 0 \\ X \\ z \end{pmatrix} = \begin{pmatrix} -\mathfrak{g}(Y, X) \\ \nabla_Y X - zY \\ z'' \end{pmatrix}$$

now the middle piece is a section of U as well, so there exists a section

$$v_1 = \begin{pmatrix} a \\ 0 \\ z'' \end{pmatrix}$$

in K , with $a \neq 0$. Since K^\perp must be orthogonal to this vector, $K^\perp \cap E^\perp$ must be of the form

$$\begin{pmatrix} 0 \\ B \\ 0 \end{pmatrix},$$

and the similar result goes for $K \cap E^\perp$. Consequently, by essentially the same argument as for Lemma 4.6, we have

Lemma 4.10. For any $A \in \Gamma(T)$, $\nabla_A X$ is a section of U .

We may, as before, choose frames (X^j) and (B^k) for the bundles U and U^\perp such that the frames commute. Then

$$\nabla_{X^j} B^k = \nabla_{B^k} X^j = 0.$$

This implies that the curvature tensor of ∇ splits into two components, its curvature on U and its curvature on U^\perp . The Ricci tensor does the same, (see Lemma 4.9), as does the rho tensor, since U and U^\perp are orthogonal. So

$$P = P_1 + P_2.$$

Here P_1 is a section of $U \otimes U$ and P_2 a section of $U^\perp \otimes U^\perp$. We shall soon see that these are rho tensors in their own right.

We now aim to prove:

Lemma 4.11. The connection ∇ is metric.

Proof. Consider the section v_1 in K , and

$$\vec{\nabla}_B v_1 = \begin{pmatrix} \nabla_B a \\ z'' B + \mathbf{ag}P(B) \\ \nabla_B z'' \end{pmatrix}.$$

The middle term $z'' B + \mathbf{ag}P(B) = z'' B + \mathbf{ag}P_2(B)$ must be zero, showing that $\mathbf{g}^{jk}(P_2)_{ij}$ is some multiple of the identity on $U^\perp[1]$ — hence that P_2 is a symmetric tensor. As the same is true of P_1 , ∇ has symmetric rho tensor, hence symmetric Ricci tensor, hence preserves a volume form, hence preserves a metric g in the conformal class. \square

Defining $h_1 = g|_U$, $h_2 = g|_{U^\perp}$, one can, exactly as in Proposition 4.8, get the proof of the decomposition:

Proposition 4.12. Locally, $(M, g) = (N_1, h_1) \times (N_2, h_2)$, where N_1 is a leaf of U and N_2 is a leaf of U^\perp .

Moreover, we have shown that P_1 and P_2 are multiples of h_1 and h_2 respectively; consequently Ric_1 and Ric_2 are as well, so both N_1 and N_2 are Einstein manifolds, with coefficients λ_1 and λ_2 . We now aim to show the relation between these coefficients.

The scalar curvature R of ∇ is $l\lambda_1 + (n - l)\lambda_2$. Hence the rho tensor, by Eq. (3), is:

$$\begin{aligned} P_1 &= -\frac{1}{n-2} \left(\text{Ric}_1 - \frac{1}{2n-2} \text{Rh}_1 \right) \\ &= -\frac{(2n-2-l)\lambda_1 + (l-n)\lambda_2}{(n-2)(2n-2)} h_1. \\ P_2 &= -\frac{1}{n-2} \left(\text{Ric}_2 - \frac{1}{2n-2} \text{Rh}_2 \right) \\ &= -\frac{(-l)\lambda_1 + (n-2+l)\lambda_2}{(n-2)(2n-2)} h_2. \end{aligned}$$

Now there is a section

$$v_1 = \begin{pmatrix} 1 \\ 0 \\ f \end{pmatrix}$$

of K (we may freely use 1, as we have established that ∇ is metric, hence got an isomorphism $L^1 \cong \mathbb{R} \times M$), and a corresponding section

$$v_2 = \begin{pmatrix} 1 \\ 0 \\ f' \end{pmatrix}$$

of K^\perp . Since v_2 is orthogonal to v_1 , $f' = -f$. Then

$$\vec{\nabla}_B v_1 = \begin{pmatrix} 0 \\ fB - P_2(B) \\ \nabla_B f \end{pmatrix}$$

as a consequence of this, we see that f is a constant and

$$f = -\frac{(2n - 2 - l)\lambda_1 + (l - n)\lambda_2}{(n - 2)(2n - 2)}.$$

carrying out a similar operation on v_2 yields the following formula

$$f = \frac{(-l)\lambda_1 + (n - 2 + l)\lambda_2}{(n - 2)(2n - 2)}.$$

Equating these terms and rearranging gives us the required

$$(n - l - 1)\lambda_1 = (1 - l)\lambda_2.$$

There is, however, a rather more fundamental reason for this seemingly arbitrary equality. For:

Proposition 4.13. *The condition*

$$(n - l - 1)\lambda_1 = (1 - l)\lambda_2$$

is equivalent to the rho tensor P_{N_1} of $\nabla|_{N_1}$ being equal to the restriction of the rho tensor on M ,

$$P_{N_1} = P|_U = P_1.$$

Proof.

$$\begin{aligned} P_1 - P_{N_1} &= \left(-\frac{(2n - 2 - l)\lambda_1 + (l - n)\lambda_2}{(n - 2)(2n - 2)} - \frac{-\lambda_1}{2(l - 1)} \right) h_1 \\ &= ((n - l - 1)\lambda_1 - (1 - l)\lambda_2) \left(\frac{(n - l)}{(l - 1)(n - 2)(2n - 2)} \right) h_1 \end{aligned}$$

Similarly

$$\begin{aligned} P_2 - P_{N_2} &= \left(-\frac{(-l)\lambda_1 + (n - 2 + l)\lambda_2}{(n - 2)(2n - 2)} - \frac{-\lambda_2}{2(n - l - 1)} \right) h_2 \\ &= ((n - l - 1)\lambda_1 - (1 - l)\lambda_2) \left(\frac{l}{(l - 1)(n - 2)(2n - 2)} \right) h_2. \end{aligned}$$

Consequently, $P_1 = P_{N_1}$ if and only if $P_2 = P_{N_2}$, and if and only if $(n - l - 1)\lambda_1 = (1 - l)\lambda_2$. \square

This is the essence of the decomposition: because of this result, $\vec{\nabla}$ operates on \mathcal{T}_{N_1} along $TN_1 = U$ just as the reduced Tractor connection $\vec{\nabla}_{N_1}$ does. Now let v_2 be the Einstein vector in \mathcal{T}_{N_1} ; then $\vec{\nabla}$ along $TN_2 = U^\perp$ will operate trivially on

$$K = v_2^\perp \cap \mathcal{T}_{N_1},$$

since K is spanned by v_1 and by sections of $(0, TN_1, 0)$. Consequently the holonomy algebra of $\vec{\nabla}$ restricted to K is $\vec{h}\vec{\sigma}|_{N_1}$.

The similar result holds for K^\perp . Thus, since $K \oplus K^\perp = \mathcal{T}$,

$$\vec{h}\vec{\sigma} = \vec{h}\vec{\sigma}|_{N_1} \oplus \vec{h}\vec{\sigma}|_{N_2}.$$

To reverse this decomposition, define (M, g) as $(N_1, h_1) \times (N_2, h_2)$ with N_1 and N_2 Einstein with Einstein coefficients related as above. Then the overall Tractor connection $\overrightarrow{\nabla}$ will be generated by $\overrightarrow{\nabla}_{N_1}$ and $\overrightarrow{\nabla}_{N_2}$ as above. Then let v_2 be the Einstein vector of \mathcal{T}_{N_1} . Then the bundle

$$K = v_2^\perp \cap \mathcal{T}_{N_1},$$

is preserved by $\overrightarrow{\nabla}$ as is its orthogonal complement

$$K = v_1^\perp \cap \mathcal{T}_{N_2},$$

where v_1 is the Einstein vector of \mathcal{T}_{N_2} . Note that $v_1 \in \Gamma(K)$ and $v_2 \in \Gamma(K^\perp)$, which explains the somewhat odd numbering of them.

Example. To illustrate these proofs, we can see that $S^4 \times \mathbb{R}^4$ does not have any holonomy-conserved subbundles in its Tractor connection (in fact it has full holonomy), while $S^4 \times \mathbb{H}^4$ is conformally flat, for \mathbb{H}^4 the hyperbolic 4-space.

Remark. Some old results of Brinkmann [7,8] can be proved directly using this decomposition theorem. For instance, the fact that any 4-manifold with two distinct Einstein structures in the conformal class is conformally flat (a direct consequence of the flatness of any Möbius structure with reduced holonomy, see next section). In our setting, the preserved subbundle spanned by the two Einstein vectors decomposes the manifold into a direct product of three- and one-dimensional Einstein spaces. But both these spaces are conformally flat, so our original manifold has trivial holonomy; in other words, it is conformally flat.

Remark. Analogously to the previous remark, we can see that if not conformally flat, a five-dimensional manifold can have up to two linearly independent Einstein structures, a six-dimensional manifold can have three, an n -dimensional manifold $n - 3$.

5. Einstein spaces: Metric cones

In this section we will give a full classification of the possible Tractor holonomies of the non-Ricci-flat Einstein spaces, using to this effect the construction of a metric cone, whose Levi-Civita holonomy corresponds to the Tractor holonomy of the original manifold.

Remark. As we mentioned in the introduction, this metric cone construction is related to the ambient metric construction of [18,15], for conformally Einstein manifolds. The actual relation is slightly subtle, the formal existence is proved in [18] and an explicit construction is given in [28]. This also provides a direct proof of a result of [21], namely that the ambient metric construction always exists if the manifold is conformally Einstein.

Definition 5.1. A conformal manifold M is said to be *indecomposable* if it cannot be decomposed into Einstein spaces as in the previous section. In other words, $\overrightarrow{\nabla}$ may preserve a single line bundle (and its orthogonal complement), but nothing else.

Remark 5.1. In the non-Ricci-flat Einstein case, indecomposable implies that the Tractor holonomy acts irreducibly on \mathbb{R}^{n+1} or $\mathbb{R}^{n,1}$ (since the only preserved line bundle is positive or negative definite).

As all Einstein manifolds of dimension 3 are conformally flat, we shall assume $n > 3$.

Let (M, g) be an Einstein manifold, $\text{Ric} = \lambda g$, $\lambda \neq 0$.

Theorem 5.2 (Einstein Classification). *The Tractor holonomy of M^n is one of the following, $n \geq 4$:*

- $SO(n, 1)$,
- $SO(n + 1)$,
- $SU(m)$ for $2m = n + 1$,
- $Sp(m)$ for $4m = n + 1$,
- G_2 for $n = 6$,
- $\text{Spin}(7)$ for $n = 7$.

Moreover, all these holonomy groups actually occur.

Remark. It is interesting to note that there is only a single holonomy possible for an indecomposable Einstein manifold with negative constant.

The remainder of this chapter will be dedicated to proving [Theorem 5.2](#).

Definition 5.3. Given an Einstein manifold (M, g) , we define the metric cone on M as $(N = \mathbb{R}^+ \times M, h)$ with

$$h = \frac{1}{\mu} dt^2 + t^2 g,$$

and $\mu = \frac{\lambda}{n-1}$.

Note that h is of definite signature if and only if M has positive Einstein constant. In the negative case, we call (N, h) a Lorentzian cone.

Then defining ∇ as the Levi-Civita connection of N , and remembering the formula:

$$2\langle \nabla_X Y, Z \rangle = X \cdot \langle Y, Z \rangle + Y \cdot \langle X, Z \rangle - Z \cdot \langle X, Y \rangle + \langle [X, Y], Z \rangle + \langle [Z, X], Y \rangle + \langle [Z, Y], X \rangle,$$

we can calculate the following equalities. For $T = \frac{\partial}{\partial t}$, and X_i a local basis of vector fields of M , extended trivially to N :

$$\begin{aligned} \nabla_T T &= 0, \\ \nabla_{X_i} T &= \frac{1}{t} X_i, \\ \nabla_T X_i &= \frac{1}{t} X_i, \\ \nabla_{X_i} X_j &= \tilde{\nabla}_{X_i} X_j - t \mu g(X_i, X_j) T, \end{aligned}$$

with $\tilde{\nabla}$ the Levi-Civita connection of g .

Given a path τ in $\{1\} \times M$, with tangent vector field Z , let Y be the parallel transport of a vector along the path, thus $\nabla_Z Y = 0$. Split Y as $Y^\perp + aT$, with $Y^\perp \in \Gamma(TM)$. Then we get the following result:

Lemma 5.4. *Extend Z and Y in the T direction, with $Z(t, x) = Z(x)$, and*

$$Y(t, x) = \frac{1}{t} Y^\perp(x) + a(x) T.$$

Then $\nabla_T Y = 0$ and $\nabla_Z Y = 0$ on $\tau \times \mathbb{R}^+$.

Proof. The function a is independent of t , so $T(a) = 0$. Hence $\nabla_T(aT) = 0$. Furthermore,

$$\begin{aligned} \nabla_T \left(\frac{1}{t} Y^\perp \right) &= T \left(\frac{1}{t} \right) Y^\perp + \frac{1}{t^2} Y^\perp \\ &= \left(-\frac{1}{t^2} + \frac{1}{t^2} \right) Y^\perp \\ &= 0, \end{aligned}$$

so $\nabla_T Y = 0$.

We can expand out the original equation $\nabla_Z Y = 0$ at $t = 1$, giving:

$$0 = \tilde{\nabla}_Z Y^\perp - \mu g(Z, Y^\perp) T + aZ + Z(a)T.$$

By linearity, this is equivalent to the two equations $0 = \tilde{\nabla}_Z Y^\perp + aZ$ and $0 = (-\mu g(Z, Y^\perp) + Z(a))T$.

Then similarly expanding $\nabla_Z Y$ for varying t :

$$\begin{aligned} \nabla_Z Y &= \frac{1}{t} \tilde{\nabla}_Z Y^\perp - \frac{t}{t} \mu g(Z, Y^\perp) T + \frac{a}{t} Z + Z(a)T \\ &= \frac{1}{t} (\tilde{\nabla}_Z Y^\perp + aZ) + (-\mu g(Z, Y^\perp) + Z(a))T \\ &= 0. \quad \square \end{aligned}$$

Now let $\check{\tau}$ be another path with same end points as τ and such that $\pi \circ \tau = \pi \circ \check{\tau}$ under the projection $\pi : N \rightarrow M$. The tangent field of $\check{\tau}$ is $\check{Z} = Z + \alpha T$, implying that $\nabla_{\check{Z}} Y = 0$. So the parallel transport of a vector along τ and $\check{\tau}$ are the same, implying that when we are computing the holonomy of ∇ , we only need to consider paths in $\{1\} \times M \cong M$.

We can now turn to the Tractor connection $\overrightarrow{\nabla}$ on \mathcal{T} , for the conformal structure $\mathbf{g} \simeq [g]$. Using the splitting given by the metric g , we can see the formal similarities with ∇ at $t = 1$.

As g is Einstein, with coefficient λ , then $\overrightarrow{\nabla} \left(\frac{n-1}{\lambda}, 0, -\frac{1}{2} \right) = 0$. Furthermore, for $R = \left(\frac{n-1}{\lambda}, 0, \frac{1}{2} \right)$, then:

$$\overrightarrow{\nabla}_{X_i} R = \begin{pmatrix} 0 \\ X_i \\ 0 \end{pmatrix} \quad \text{and} \quad \overrightarrow{\nabla}_{X_i} \begin{pmatrix} 0 \\ X_j \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ \tilde{\nabla}_{X_i} X_j \\ 0 \end{pmatrix} - \mu g(X_i, X_j) R.$$

Hence under the formal identification of R with T and $(0, X_i, 0)$ with X_i , we get $\nabla_Z \cong \overrightarrow{\nabla}_Z$ for $Z \in \Gamma(T)$ at $t = 1$. Then by the previous lemma and its implication for the holonomy of ∇ , the next theorem is proved:

Theorem 5.5. *The holonomy groups of $(\mathcal{T}, \overrightarrow{\nabla}, (M, \mathbf{g}))$ and $(TN, \nabla, (N, h))$ are isomorphic.*

Hence the holonomy of $\overrightarrow{\nabla}$ is metric, and irreducible by Remark 5.1, and must be one of those classified by Merkulov and Schwachhöfer in [30]. In the negative Einstein case, a look at the table shows that the only possible holonomy is the full $SO(n, 1)$ group itself. For the positive Einstein, we need the following result:

Proposition 5.6. *The metric cone (N, h) is Ricci-flat.*

Proof. By the definition of h and the corresponding ∇ , the curvature $R = R^\nabla$ is

$$R_{T,-} = 0, \\ R_{X_i X_j X_k} = \tilde{R}_{X_i X_j X_k} - \mu g(X_j, X_k) X_i + \mu g(X_i, X_k) X_j,$$

with \tilde{R} the curvature of $\tilde{\nabla}$.

Then taking traces,

$$\text{Ric}(T, -) = 0, \\ \text{Ric}(X_j, X_k) = \tilde{\text{Ric}}(X_j, X_k) + (1 - n)\mu g(X_j, X_k) \\ = \lambda g(X_j, X_k) - \lambda g(X_j, X_k) \\ = 0. \quad \square$$

So the possible holonomies reduce to those corresponding to metrics which are Ricci-flat, namely $SO(n + 1)$, $SU(m)$, $Sp(m)$, G_2 and $\text{Spin}(7)$.

The $SO(n + 1)$ case is generic. The $SU(m)$ holonomy on the cone corresponds to Sasaki–Einstein manifolds, the $Sp(m)$ to 3-Sasakian ones, G_2 and $\text{Spin}(7)$ to weak holonomy manifolds [5]; all of which can be realised on cones over compact manifolds, implying that all exist as holonomy groups of conformal Tractor connections on compact manifolds.

It is immediate that a metric cone on any one-dimensional space is flat. We now aim to show that the metric cone on a two-dimensional Einstein space of constant scalar curvature is also flat.

Proposition 5.7. *Any Tractor connection in two dimensions with a preserved Tractor u is flat.*

Proof. In this case, using the metric defined by $\pi^2(u)$ to identify T and T^* , we have,

$$P = \mu Id,$$

With μ a constant. However, [11], the only curvature element of a Tractor/Möbius connection in two dimensions is the Cotton–York tensor — which must vanish entirely, as $\nabla P = 0$, making the connection flat. \square

Remark. In [6] Baston presents a local twistor theory, which, in the case of conformal manifolds, is just given by the vector bundle associated to \mathcal{G}_0 under the spin representation of G_0 , and the extension of $\overrightarrow{\nabla}$ to this new context.

A parallel section of this bundle is equivalent, by [19], with the existence of a spinor ψ on the underlying manifold solving the *twistor equation* for all vector fields X :

$$\tilde{\nabla}_X \psi + \frac{1}{n} X \cdot D\psi = 0,$$

with D the Dirac operator. Paper [23] by Habermann analyses solutions to this twistor equation; she shows that these imply that the manifold is conformally Einstein, of non-negative scalar curvature. So twistor–spinors should be visible in the cone construction.

A connection preserves an element if and only if its holonomy group does so as well. Now, since $\tilde{\nabla}$ has the same holonomy as the cone connection ∇ , it preserves a Tractor spinor if and only if ∇ preserves a spinor $\hat{\psi}$ on the cone. That equation is $0 = \nabla_X \hat{\psi} = \tilde{\nabla}_X \hat{\psi} + (1 + \frac{1}{n} \text{trace}(gP))X \cdot \hat{\psi}$. Decomposing the spinor representation in terms of irreducible representations of $\tilde{\nabla}$ gives the twistor equation.

The holonomy groups G_2 and $\text{Spin}(7)$ correspond to the existence of a preserved spinor on the cone — thus to the existence of twistor–spinors on the manifold itself.

Remark. The concept of a twistor–spinor is a generalisation of that of a Killing spinor. A Killing spinor is a spinor ψ solving the equation

$$\nabla_X \psi = \lambda X \cdot \psi$$

for all vector fields X and some constant λ . In [5], Bär showed that having a Killing spinor is equivalent with having a parallel spinor on the metric cone. So the cases of weak holonomy $SU(3)$ and nearly Kählerian structures are covered by the Tractor connection; in fact in his paper [9] constructing manifolds of exceptional holonomy, R.L. Bryant produces manifolds of holonomy G_2 and $\text{Spin}(7)$ as metric cones on $SU(3)/T^2$ and $SO(5)/SO(3)$ respectively. Thus all the holonomies listed actually occur.

We can now turn to the Ricci-flat case, which is actually simpler than the general Einstein case, but with an added subtlety.

6. Ricci-flat spaces

Let (M^n, g) be a Ricci-flat space of indecomposable Tractor holonomy. As M^n is Ricci-flat, its Tractor holonomy is contained within $SO(n) \times \mathbb{R}^n$. Fix a point $b \in M$ for calculating the holonomy groups, and let H be the metric holonomy of M , D its Tractor holonomy, $\mathfrak{h}, \mathfrak{l}$ their Lie algebras.

Then:

Lemma 6.1. $H \subset D$, or, equivalently, $\mathfrak{h} \subset \mathfrak{l}$.

Proof. Let Y be the parallel transport of a vector along a path τ with tangent field X ; in other words $\nabla_X Y = 0$, for ∇ the metric connection on M . Then

$$\tilde{\nabla}_X \begin{pmatrix} x \\ Y \\ 0 \end{pmatrix} = \begin{pmatrix} \nabla_X x - g(X, Y) \\ \nabla_X Y = 0 \\ 0 \end{pmatrix},$$

which is zero for $x = \int_\tau g(X, Y)$, proving that every metric holonomy element is a Tractor holonomy element. This argument also works in reverse, showing that $\pi(\mathfrak{l}) = \mathfrak{h}$, where π is the projection of $\mathfrak{co}(n)_b \oplus T_b$ onto its first component. \square

This demonstrates that $\mathfrak{l} \subset \mathfrak{b} = \mathfrak{h} \oplus T_b$. But first:

Lemma 6.2. *The representation of \mathfrak{h} on T_b is irreducible.*

Proof. If a bundle $S \subset T$ is preserved by ∇ , then the bundle $\mathcal{E}[1] \oplus S$ is preserved by $\tilde{\nabla}$. Thus, since we assume our Tractor holonomy to be indecomposable, \mathfrak{h} must act irreducibly on T_b . \square

Then since the Lie bracket on T_b is trivial, the adjoint representation of $\mathfrak{h} \subset \mathfrak{l} \subset \mathfrak{b}$ on the second component of \mathfrak{b} is the usual, irreducible one. Accordingly this adjoint representation splits \mathfrak{b} into two irreducible representations, isomorphic to \mathfrak{h} and T_b .

As a consequence, $\mathfrak{l} \cong \mathfrak{h}$ or $\mathfrak{l} \cong \mathfrak{b}$. We now claim that

Lemma 6.3. $\mathfrak{l} \cong \mathfrak{b}$.

Proof. Reasoning by contradiction, we assume that $\mathfrak{l} \cong \mathfrak{h}$, and go on to show that this violates our indecomposability assumption.

Express \mathfrak{b} as $\mathfrak{h}_0 \oplus \mathbb{R}_0^n$, the sum of the irreducible representations of \mathfrak{h} . Then, as $\mathfrak{h}_0 \cong \mathfrak{h}$ acts irreducibly on T_b , there is, at b , a new splitting of \mathcal{T} corresponding to the splitting $\mathbb{R}^{n*} \oplus \mathfrak{h}_0 \oplus \mathbb{R}_0^n$. This splitting is

$$\mathcal{E}[-1]_0 \oplus T[-1]_0 \oplus \mathcal{E}[-1].$$

Then $\mathfrak{h} \subset \mathfrak{l}$ preserves the new vectors $(1, 0, 0)$ and $(0, 0, 1)$. This shows that \mathfrak{l} preserves a rank 2 subbundle, contradicting indecomposability. \square

Putting this together, we can now claim the following theorem:

Theorem 6.4. *The possible indecomposable Tractor holonomy groups for the conformal manifold (M, \mathbf{g}) , conformally Ricci-flat, are:*

- $SO(n) \times \mathbb{R}^n, n \geq 4$ (generic),
- $SU(m) \times \mathbb{R}^{2m}, m \geq 2$ (Calabi–Yau),
- $Sp(m) \times \mathbb{R}^{4m}, m \geq 1$ (hyper-Kähler),
- $G_2 \times \mathbb{R}^7$ (see [9]),
- $Spin(7) \times \mathbb{R}^8$ (see [9]),

and all of these groups do occur.

Remark. This result offers an alternative proof for the theorems in Listing’s paper [29], in the special case of conformally Ricci-flat manifolds.

Remark. The metric cones constructed in the previous chapter are Ricci-flat (pseudo-)Riemannian manifolds. They are not, however, indecomposable; in fact, their Tractor holonomy is equal to their metric holonomy. This property characterises metric cones.

7. Addendum: Symmetric spaces

A symmetric space (S, g) is a manifold such that $\nabla^g R^g = 0$ for R^g the full curvature tensor. Since symmetric spaces are direct products of Einstein spaces, $\nabla^g \mathbf{P} = 0$, so $\nabla^g W^g = 0$ as well, for the Weyl tensor W^g .

Kobayashi and Nomizu [24] define the infinitesimal holonomy of a connection $\overrightarrow{\nabla}$ as the span of the iterated derivatives of its curvature tensor

$$\overrightarrow{\nabla}_{X_1} \overrightarrow{\nabla}_{X_2} \dots \overrightarrow{\nabla}_{X_{m-2}} \overrightarrow{R}(X_{m-1}, X_m),$$

estimated on any vector fields $X_1 \dots X_m$. They demonstrate that the infinitesimal holonomy is contained in the holonomy algebra.

The curvature of the Tractor connection is $\overrightarrow{R} = W^g$. Now $\overrightarrow{\nabla}_X = X + \nabla_X^g + \mathbf{P}(X)$ and the middle element has no effect on either W^g or on \mathbf{P} (since S is a direct product of Einstein spaces). Consequently if we look at the span of the iterated derivative of W^g , we are looking at the algebraic span of W^g under the action of $X + \mathbf{P}(X)$ for all X . Thus the infinitesimal holonomy at any point $p \in S$ is a representation of the algebra generated by $X_p + \mathbf{P}(X)_p$. And if \mathbf{P} is not identically zero, this algebra is isomorphic with one of $\mathfrak{so}(n + 1, 1)$, $\mathfrak{so}(n + 1)$ or $\mathfrak{so}(n, 1)$.

This demonstrates that any manifold that is conformal to a symmetric space is either conformally flat or has the maximal holonomy in its category — $SO(n + 1, 1)$ if the symmetric space is not Einstein, and $SO(n + 1)$ or $SO(n, 1)$ if it is (Ricci-flat symmetric spaces are flat).

Example. These results give an independent proof to the results of Leitner [27], that the conformal holonomy of $SO(4)$, locally isomorphic to $S^3 \times S^3$, is $SO(7)$. The group $SO(4)$ is a positive Einstein symmetric space, not conformally flat (consider the fate of the Tractor vector $(1, 0, 0)$ under parallel translation) so the result follows.

Example. Note that the same argument shows that the manifold $S^n(a) \times S^n(b)$ where $a \neq b$ are the radii of the spheres, has full holonomy $SO(2n + 1, 1)$.

Two similar results also are implied:

Example. If a manifold (M, \mathbf{g}) is conformal to an Einstein symmetric space, then it cannot be conformal to any other Einstein space, or any other symmetric space, unless it is conformally flat.

And:

Example. If (M, \mathbf{g}) is conformal to a symmetric space in two different ways, then its tractor holonomy is full or null.

8. Indefinite signature

Most of the results of this paper extend to the general pseudo-Riemannian case, see [3]. The cone construction and the Ricci-flat results still apply, as do the results about umbilicity. The decomposition theorem, however, requires an extra condition: that $K \cap K^\perp$ be of rank 1 or 0 (a condition that is automatically true in the definite signature case). This is equivalent with requiring that $U \cap U^\perp = 0$; without it, the decomposition cannot proceed.

And, unlike the result for $\mathfrak{so}(n + 1, 1)$ proved in [17], there are non-trivial subalgebras of $\mathfrak{so}(p + 1, q + 1)$ acting irreducibly on $\mathbb{R}^{(p+1, q+1)}$. Thus we have many other candidate algebras to deal with.

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